

## The *GALEX* Arecibo SDSS Survey – II. The star formation efficiency of massive galaxies

David Schiminovich,<sup>1\*</sup> Barbara Catinella,<sup>2</sup> Guinevere Kauffmann,<sup>2</sup> Silvia Fabello,<sup>2</sup> Jing Wang,<sup>2,3</sup> Cameron Hummels,<sup>1</sup> Jenna Lemonias,<sup>1</sup> Sean M. Moran,<sup>4</sup> Ronin Wu,<sup>5</sup> Riccardo Giovanelli,<sup>6</sup> Martha P. Haynes,<sup>6</sup> Timothy M. Heckman,<sup>4</sup> Antara R. Basu-Zych,<sup>7</sup> Michael R. Blanton,<sup>5</sup> Jarle Brinchmann,<sup>8,9</sup> Tamás Budavári,<sup>4</sup> Thiago Gonçalves,<sup>10</sup> Benjamin D. Johnson,<sup>11</sup> Robert C. Kennicutt,<sup>11,12</sup> Barry F. Madore,<sup>13</sup> Christopher D. Martin,<sup>10</sup> Michael R. Rich,<sup>14</sup> Linda J. Tacconi,<sup>15</sup> David A. Thilker,<sup>4</sup> Vivienne Wild<sup>16</sup> and Ted K. Wyder<sup>10</sup>

<sup>1</sup>Department of Astronomy, Columbia University, New York, NY 10027, USA

<sup>2</sup>Max-Planck Institut für Astrophysik, D-85741 Garching, Germany

<sup>3</sup>Center for Astrophysics, University of Science and Technology of China, 230026 Hefei, China

<sup>4</sup>Department of Physics and Astronomy, The Johns Hopkins University, Baltimore, MD 21218, USA

<sup>5</sup>Department of Physics, New York University, New York, NY 10003, USA

<sup>6</sup>Center for Radiophysics and Space Research, Cornell University, Ithaca, NY 14853, USA

<sup>7</sup>NASA Goddard Space Flight Center, Laboratory for X-ray Astrophysics, Greenbelt, MD 20771, USA

<sup>8</sup>Leiden Observatory, Leiden University, 2300 RA, Leiden, the Netherlands

<sup>9</sup>Centro de Astrofísica, Universidade do Porto, 4150-762 Porto, Portugal

<sup>10</sup>California Institute of Technology, Pasadena, CA 91125, USA

<sup>11</sup>Institute of Astronomy, Cambridge CB3 0HA

<sup>12</sup>Steward Observatory, University of Arizona, Tucson, AZ 85721, USA

<sup>13</sup>Observatories of the Carnegie Institution of Washington, Pasadena, CA 91101, USA

<sup>14</sup>Department of Physics and Astronomy, University of California, Los Angeles, CA 90095, USA

<sup>15</sup>Max Planck Institut für extraterrestrische Physik, D-85741 Garching, Germany

<sup>16</sup>Institut d'Astrophysique de Paris, 75014 Paris, France

Accepted 2010 June 15. Received 2010 May 27; in original form 2010 March 18

### ABSTRACT

We use measurements of the H I content, stellar mass and star formation rates (SFRs) in  $\sim 190$  massive galaxies with  $M_{\star} > 10^{10} M_{\odot}$ , obtained from the *GALEX* (*Galaxy Evolution Explorer*) Arecibo SDSS (Sloan Digital Sky Survey) survey described in Paper I to explore the global scaling relations associated with the bin-averaged ratio of the SFR over the H I mass (i.e.  $\Sigma \text{SFR} / \Sigma M_{\text{H I}}$ ), which we call the H I-based star formation efficiency (SFE). Unlike the mean specific star formation rate (sSFR), which decreases with stellar mass and stellar mass surface density, the SFE remains relatively constant across the sample with a value close to  $\text{SFE} = 10^{-9.5} \text{ yr}^{-1}$  (or an equivalent gas consumption time-scale of  $\sim 3 \times 10^9 \text{ yr}$ ). Specifically, we find little variation in SFE with stellar mass, stellar mass surface density,  $\text{NUV} - r$  colour and concentration ( $R_{90}/R_{50}$ ). We interpret these results as an indication that external processes or feedback mechanisms that control the gas supply are important for regulating star formation in massive galaxies. An investigation into the detailed distribution of SFEs reveals that approximately 5 per cent of the sample shows high efficiencies with  $\text{SFE} > 10^{-9} \text{ yr}^{-1}$ , and we suggest that this is very likely due to a deficiency of cold gas rather than an excess SFR. Conversely, we also find a similar fraction of galaxies that appear to be gas-rich

\*E-mail: ds@astro.columbia.edu

for their given sSFR, although these galaxies show both a higher than average gas fraction and lower than average sSFR. Both of these populations are plausible candidates for ‘transition’ galaxies, showing potential for a change (either decrease or increase) in their sSFR in the near future. We also find that  $36 \pm 5$  per cent of the total H I mass density and  $47 \pm 5$  per cent of the total SFR density are found in galaxies with  $M_{\star} > 10^{10} M_{\odot}$ .

**Key words:** galaxies: evolution – galaxies: fundamental parameters – radio lines: galaxies – ultraviolet: galaxies.

## 1 INTRODUCTION

Measurements of volume-averaged stellar mass, star formation rate (SFR) and gas densities over cosmic time provide fundamental constraints on galaxy evolution by describing the integrated past history, present activity and future potential for galaxy growth. In the past decade, substantial progress has been made in accurately measuring the first two – the stellar mass function (Bell et al. 2003; Borch et al. 2006) and SFR density (Brinchmann et al. 2004; Salim et al. 2007; Wyder et al. 2007). On the other hand, quantifying gas content across redshift, and in particular identifying the gas supply and reservoir most closely linked to recent and future star formation (SF), has remained challenging.

Recent progress has been made in the local Universe through blind H I surveys such as HI Parkes All-Sky Survey (HIPASS; Meyer et al. 2004) and Arecibo Legacy Fast ALFA (ALFALFA; Giovanelli et al. 2005) which are yielding H I mass functions of representative volumes (Zwaan et al. 2005; Stierwalt et al. 2009) and optical identifications of nearly all of H I-detected sources (e.g. Saintonge et al. 2008; Martin et al. 2009). This latter result, the fact that blind H I surveys do not reveal a population of H I-rich galaxies without optical counterparts, suggests that the study of the gas content of optically selected samples can provide a nearly complete census of H I in the local Universe.

In this paper, rather than performing a complete census of H I, we focus exclusively on a subsample of massive galaxies in the local Universe that have been homogeneously observed by the Sloan Digital Sky Survey (SDSS), *Galaxy Evolution Explorer* (GALEX) and now Arecibo as part of the GASS survey (described in Catinella et al. 2010, hereafter Paper I). Briefly stated, the goals of the GASS survey are to measure the H I properties of  $\sim 1000$  massive ( $M_{\star} > 10^{10} M_{\odot}$ ) galaxies in the local Universe ( $z < 0.05$ ), in a mass range that includes the transition mass above which most galaxies have ceased forming stars. To maximize survey efficiency, galaxies are observed to a gas fraction limit of 2–5 per cent, ensuring that the quantity of gas probed remains relevant for on-going and future SF. In Paper I, we found that the H I gas fraction decreased with stellar mass, stellar mass surface density and NUV –  $r$  colour.

Observations and models suggest that SF in massive galaxies may be quenched by internal [active galactic nuclei (AGN) or SF feedback] or external means (stripping or other environmental effects). This has been used to explain both the decline of the global SFR and the growth of stellar mass on the red sequence. If quenching is due to the depletion of a gaseous reservoir, one might expect its signature to be evident in a decrease in the mean H I gas fraction. However, if the quenching is due to internal processes that inhibit SF (e.g. Martig et al. 2009) then the signature of quenching might also be reflected in a decreased star formation efficiency (e.g.  $SFE = SFR/M_{\text{H}}$ ) or the equivalent converse, an increased gas consumption time-scale ( $t_{\text{cons}} = M_{\text{H}}/SFR$ ). While SFR efficiencies have been discussed extensively (e.g. Roberts 1963; Larson, Tinsley &

Caldwell 1980; Kennicutt 1983; Kennicutt, Tamblyn & Congdon 1994; Boselli et al. 2001; Bothwell, Kennicutt & Lee 2009), this work marks the first time that it has been determined for a representative sample of galaxies selected exclusively by stellar mass.

In this paper, we measure the distribution of SFR efficiencies across the GASS sample and ask whether we find an excess of highly efficient or inefficient star-forming galaxies? We also investigate how the mean efficiency varies across the sample. Additionally, we take advantage of our simple selection criterion to produce a determination of several volume-averaged physical quantities. We combine GASS and GALEX measurements of  $\sim 190$  galaxies with recent determinations of the local stellar mass function to calculate the H I mass density and SFR density as a function of stellar mass. This analysis allows us to compare the properties of our mass-selected sample to that of the full population. We emphasize here that the sample, while restricted to relatively massive galaxies, includes many galaxies that are blue and/or not quiescent, in contrast with samples selected by colour or early-type morphology, the latter of which are known to have low gas fractions and low SFRs (e.g. Bregman, Hogg & Roberts 1992; Yi et al. 2005; Morganti et al. 2006; Oosterloo et al. 2007). In fact, as we discuss below, a significant fraction of the total SFR in the local Universe is taking place within galaxies in this stellar mass range.

We use these measurements to address several fundamental questions: how much of the total H I in the local Universe is associated with massive galaxies? How does this change across the so-called transition mass of  $10^{10.5} M_{\odot}$ ? How does the H I density compare with the measured SFR density over the same mass range, and what does this suggest regarding the global SFR efficiency of massive galaxies? What fraction of this gas is in the process of efficiently forming stars, versus the fraction that may be building up a reservoir for future SF? We interpret our results in the context of quenching models, as well as scenarios that may lead to the return of a galaxy back on to the star-forming sequence.

Throughout this paper, we make use of the flat  $\Lambda$  cold dark matter cosmology with  $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$  and  $\Omega_{\Lambda} = 0.7$ .

## 2 DATA

### 2.1 H I, optical data and derived quantities

Our H I data are taken from the GASS first data release (DR1) described in Paper I, and we refer the reader there for details on the observations and initial data analysis. We briefly review the most pertinent information. Our large parent sample (PS) contains 12 006 galaxies with  $M_{\star} > 10^{10} M_{\odot}$  and  $0.025 < z < 0.05$  visible from Arecibo and located within the footprint of the SDSS primary spectroscopic survey, the projected GALEX Medium Imaging Survey and ALFALFA. The first release of GASS data contains  $\sim 190$  galaxies of which 176 are new measurements, with 99 detections and 77 upper limits. Another  $\sim 10$ –15 previously detected gas-rich

galaxies (from ALFALFA and other surveys) are added to produce the statistically representative, volume-limited sample (DR1) used in this paper. DR1 contains 20 per cent of the galaxies planned for the full GASS survey.

H I masses are calculated for the detected galaxies, and upper limits determined for the non-detections. For this measurement, our targets are considered to be unresolved by the Arecibo 3.5 arcmin beam. Additionally, no correction is made for self-absorption. Detected H I masses range from  $4.6 \times 10^8$  to  $3.2 \times 10^{10} M_\odot$ . Upper limits ( $5\sigma$ ) from non-detections assume a  $300 \text{ km s}^{-1}$  velocity width. Upper limits are indicated on figures using arrows plotted at the location of the  $5\sigma$  detection limit. Although velocity widths have been measured from the Arecibo spectrum, we do not make use of those quantities in this paper. We have investigated possible source confusion and contamination due to signal from galaxies close to or at the same redshift as our target galaxies. While such confusion does not have an influence on our main results, where appropriate we have indicated those galaxies for which confusion might cause an overestimate of the object's H I mass.

Although the GASS PS was defined using the SDSS DR6, measurements and derived physical properties reported here were obtained using the SDSS DR7, including the Max Planck Institute for Astrophysics (MPA)/John Hopkins University (JHU) value-added catalogues. Ultraviolet (UV) photometric measurements were calculated directly from pipeline-processed *GALEX* images based on *GALEX* data release GR45. These quantities have been tabulated in Paper I. All photometric quantities used in this paper have been corrected for Galactic extinction. Stellar masses have been calculated using the SDSS photometry only using the methodology described in Salim et al. (2007). Following that work, we assume a Chabrier (2003) initial mass function (IMF) for all derived quantities based on stellar masses and SFRs, including gas mass fractions and SFR efficiencies. Stellar mass surface densities are calculated assuming that 50 per cent of the stellar mass is contained within the  $r$ -band half-light radius.

## 2.2 Star formation rates and star formation efficiencies

In this section, we describe the method we use to calculate SFRs and SFEs. SFRs in this paper are derived from UV luminosities corrected for internal dust attenuation. A particular challenge is that the GASS sample contains many galaxies with low-level SF activity for which old and/or non-star-forming components (evolved stars, AGN) can produce a UV luminosity comparable to those of a faint young stellar population. This issue is not unique to the UV – deriving SFRs for nearly passive galaxies is challenging using any SF diagnostic [e.g. UV, infrared (IR), H $\alpha$ , radio continuum, etc.]. Furthermore, measuring dust attenuation in individual galaxies is a complex problem, and nearly unfeasible with low signal-to-noise ratio (S/N) data. A full treatment of these issues is beyond the scope of this paper, and we adopt a relatively simple UV-optical-based approach following analyses in previous work. Interesting alternative multiwavelength methods for deriving SFRs across the galaxy population have been developed (Salim et al. 2007; Cortese et al. 2008; Kennicutt et al. 2009), and we will explore these in future work.

A simple, one-component expression for calculating SFRs using observed UV luminosities is given by

$$\text{SFR} = \frac{L_{\text{UV}} f_{\text{UV}}(\text{young}) 10^{0.4 A_{\text{UV}}}}{\eta_{\text{UV}}},$$

where  $\eta_{\text{UV}}$  is the (SF history and metallicity-dependent) conversion factor between UV luminosity and recent-past-averaged SFR,  $A_{\text{UV}}$  is the (geometrically averaged) ratio of intrinsic UV luminosity to measured UV luminosity and  $f_{\text{UV}}(\text{young})$  is a measure of the fraction of UV light that originates in a young – as opposed to a highly evolved stellar population. All of these quantities are highly simplified and cannot be considered to be independent, but for our derivations and discussion below they will be treated as such.

Since the GASS sample is drawn from the SDSS and therefore has seven band photometry (*fnugriz*) and spectral line indices, we use dust attenuations ( $A_{\text{NUV}}$ ) derived using the methodology described in Johnson et al. (2007), slightly modified to produce more accurate SFRs for galaxies with evolved stellar populations. Variations on this methodology have been proposed (e.g. Cortese et al. 2008), and we have checked this method against ours to ensure that applying different corrections has little impact on our conclusions. We do not employ the ‘hybrid’ attenuation correction used in Schiminovich et al. (2007) because of the unavailability of the  $z$ -band dust attenuation measure ( $A_z$ ), although we have calculated similar dust attenuation measures using  $\tau_V$  in order to verify the robustness of the results reported here. We also do not make use of H $\alpha$ -derived SFRs because emission lines with sufficient S/N are only available for a subset of the sample and fibre aperture corrections can be large for our (relatively) low-redshift sample.

Johnson et al. (2006, 2007) used the IR/UV flux ratio ( $\text{IRX}$ ) from a sample of 1000 SDSS galaxies to derive a UV-based dust attenuation measure ( $A_{\text{IRX}}$ ). By combining UV-optical colours with  $D_n(4000)$ , which correlates with SF history, Johnson et al. (2007) showed that

$$A_{\text{IRX}} = 1.25 - 1.33x + 1.19y - 1.02xy,$$

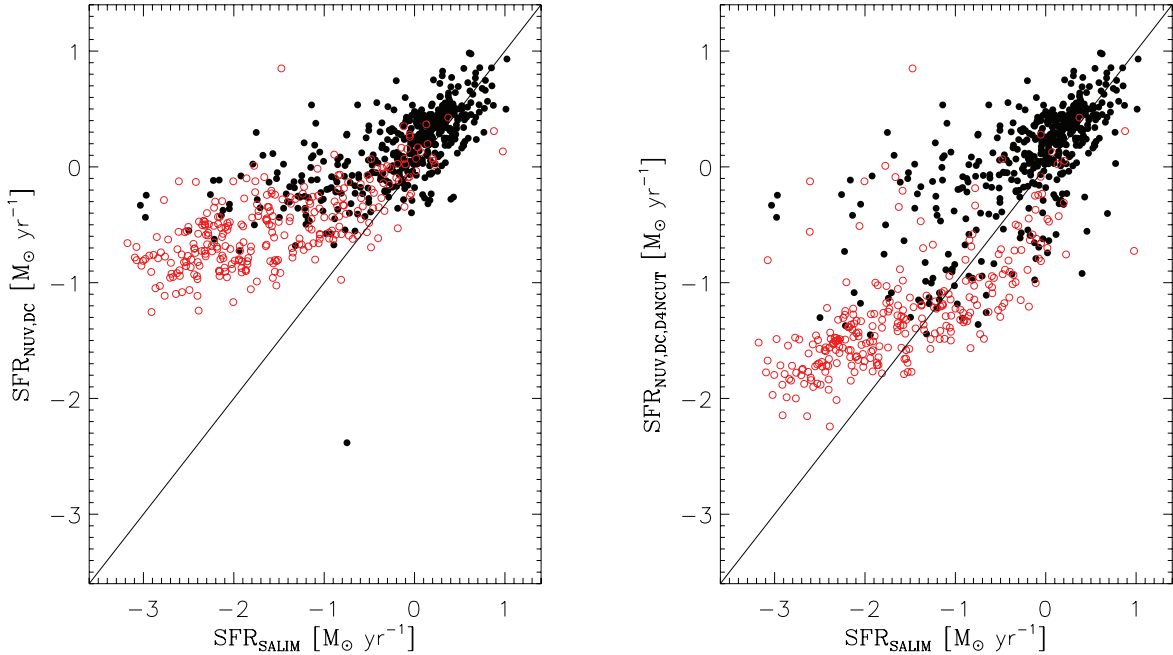
where  $x = D_n(4000) - 1.25$  and  $y = 0.1(\text{NUV} - r) - 2$  and coefficients have been taken from table 2 in Johnson et al. (2007). Johnson et al. (2007) show that the empirically derived  $A_{\text{IRX}}$  shows a close correspondence to  $A_{\text{UV}}$  for galaxies with  $D_n(4000) < 1.7$  and using the Calzetti et al. (2000) attenuation curve one can show that  $A_{\text{NUV}} = 0.81 A_{\text{IRX}}$ .

For galaxies with  $D_n(4000) > 1.7$ , the IR/UV flux ratio is also sensitive to dust heating by light from evolved stars and is not a good measure of the dust attenuation of a young stellar population. We attempt to bound the range of possible dust attenuations by considering two distinct cases. In the first (DC, D4NCUT), we assume that the IR luminosity in galaxies with  $D_n(4000) > 1.7$  is reprocessed light from an old population and therefore set  $A_{\text{NUV}} = 0$  for galaxies with  $D_n(4000) > 1.7$ . In the second case (DC), we apply our empirically derived dust attenuation corrections to the entire sample.

We make a further simplification by assuming that  $f_{\text{UV}}(\text{young}) = 1$  or that all of the UV light in our measurements comes from a young stellar population. Measurements of the UV-optical colours of early-type galaxies (e.g. Rich et al. 2005; Yi et al. 2005; Donas et al. 2007) show  $\text{NUV} - r \sim 5$ –6. Such a colour, converted into a specific star formation rate (sSFR), leads to a value close to  $\text{SFR}/M_\star \sim 10^{-12} \text{ yr}^{-1}$ . This implies that sSFRs close to this value will very likely include a contribution from old stars and will most likely lead to an overestimate of SFR, since  $f_{\text{UV}}(\text{young})$  may be much less than 1. While our assumption leads to an overestimate of SFRs in weakly star-forming galaxies, it ends up having a small effect on average quantities discussed in the first part of this paper.

SFRs were calculated assuming a constant  $\eta_{\text{NUV}} = 10^{28.165}$  and therefore

$$\text{SFR}(M_\odot \text{ yr}^{-1}) = 10^{-28.165} L_{\nu, \text{NUV}} (\text{erg s}^{-1} \text{ Hz}^{-1}) 10^{0.4 A_{\text{NUV}}}$$



**Figure 1.** Comparison of SFR derived using methodology in this paper and SFR from Salim et al. (2007), illustrating how SFRs bracket a range of possible values. In both plots, black filled points are galaxies with the SDSS  $D_n(4000) < 1.7$  and red open circles with  $D_n(4000) > 1.7$ . Left-hand panel shows comparison of dust-corrected SFR for full range of galaxies (DC sample), based on Johnson et al. (2007). Right-hand panel shows comparison with dust-corrected SFR with a  $D_n(4000)$  cut on attenuation (DC,D4NCUT sample), meaning that  $A_{\text{NUV}} = 0$  for galaxies with  $D_n(4000) > 1.7$ .

as derived by Salim et al. (2007) assuming a Chabrier (2003) IMF and a continuous recent (100–300 Myr) SF history, which makes these SFRs directly comparable to those in that work and most other recent determinations. For a standard Salpeter IMF (between 0.1 and 100  $M_\odot$ ), SFRs would be a factor of  $\sim 1.5$  higher.

In this paper, we are primarily focused on quantities derived using the ratio of mean values. In several cases, but primarily for the SFRs, these mean quantities can be dominated by a few very high SFR galaxies. Therefore, we took some care to ensure that our results were not affected by possible errors in measured SFRs for galaxies with the highest SFR. Many of these galaxies are IR luminous and detected by *IRAS*. We have compared our derived SFR and the total infrared luminosity, obtained from the IIFSCz (Imperial IRAS FSC redshift catalogue; Wang & Rowan-Robinson 2009). Overall, the agreement is sufficiently good that differences are unlikely to affect our results. It is possible, however, that we could be underestimating the SFR for the most IR-luminous galaxies, thereby underestimating the SFEs for a subset of our sample.

In Fig. 1, we compare SFRs obtained for both cases described above, to a subset of GASS PS galaxies with SFRs calculated previously by Salim et al. (2007). We note that the DC,D4NCUT sample compares quite favourably to Salim et al. (2007) and adopt this as the primary sample for this paper.

Finally,  $H\text{I}$  masses, stellar masses and SFRs are used to derive  $H\text{I}$ -based gas fractions ( $M_{\text{H I}}/M_\star$ ), star formation efficiencies ( $\text{SFE} = \text{SFR}/M_{\text{H I}}$ ) and specific star formation rates ( $\text{SFR}/M_\star$ ). We follow Leroy et al. (2008) in adopting SFE, the SFR-to-gas ratio, as our primary derived quantity, as opposed to its reciprocal, often called the gas consumption time-scale (or the *Roberts time*). Although our definition is straightforward, it is not necessarily identical to the SFE term sometimes used in SF laws and related prescriptions.

For all of our derived quantities we have opted for the simplest possible definition. We do not include a gas recycling term when calculating SFE. Although helium adds 26 per cent to the cold

gas mass we do not include this factor to calculate a cold gas mass nor do we include a correction for the unmeasured molecular component. This is in contrast, for example, to the factor of 2.3 applied to  $H\text{I}$  masses in the SINGG survey (Hanish et al. 2006; Meurer et al. 2006). We do note that the molecular phase in  $\sim 1/3$  the GASS sample is currently being studied in a corollary CO survey being carried out at IRAM (COLDGASS; Saintonge et al., in preparation).<sup>1</sup>

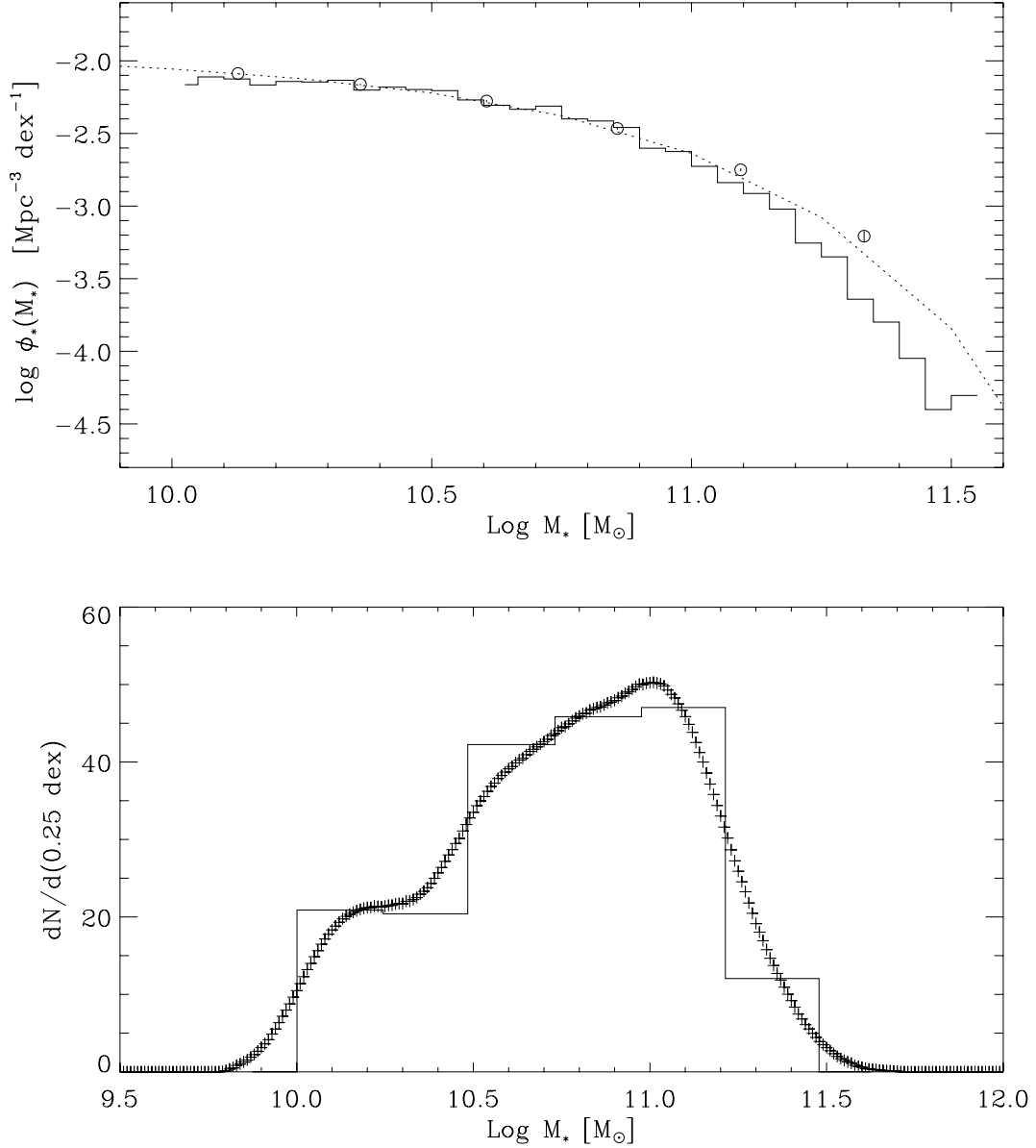
### 3 GLOBAL $H\text{I}$ , SFR AND SFE DENSITY

#### 3.1 Methodology

The very simple selection criterion used for the GASS PS allows us to combine our DR1  $H\text{I}$  measurements with the local stellar mass function to determine the volume-averaged  $H\text{I}$  mass density for massive ( $M_\star > 10^{10} M_\odot$ ) galaxies in the local Universe. Using SFR measurements we can also determine the analogous SFR density and volume-averaged SFE. These volume-weighted quantities are derived using the local stellar mass function  $\phi_\star(M_\star)$  taken from Borch et al. (2006), under the assumption that GASS DR1 observations sample an unbiased, volume-limited distribution of galaxies in any particular stellar mass bin.

We first derive weights that allow us to scale the GASS sample to match the stellar mass function of Borch et al. (2006). We also show that our weighting scheme allows us to (trivially) recover  $\phi_\star(M_\star)$  using the GASS DR1 sample. To avoid effects that might arise from binning, we express the distribution of stellar masses in the GASS DR1 as a continuous (Gaussian-smoothed) density function  $d(M_\star)$ , shown in the bottom panel of Fig. 2.

<sup>1</sup> [http://www.mpa-garching.mpg.de/COLD\\_GASS](http://www.mpa-garching.mpg.de/COLD_GASS)



**Figure 2.** Stellar mass function and distribution of GASS sample galaxies. Top panel: circles show values calculated in 0.25 dex bins using the methodology described in the paper with bootstrap errors. Solid histogram shows distribution of GASS PS galaxies selected from the SDSS DR6, with normalization scaled to fit stellar mass function. Dotted curve is stellar mass function from Borch et al. (2006). Bottom panel: histogram shows binned stellar mass distribution from the GASS DR1 sample. Plus signs show Gaussian-smoothed density function versus  $M_*$  for DR1.

For all  $i$  such that  $|\log M_{*,i} - \log M_{*,\text{bin}}| < \Delta \log M_{*,\text{bin}}/2$ :

$$\phi_*(\bar{M}_{*,\text{bin}}) = \frac{\sum_i \phi_*(M_{*,i}) w_i}{\sum_i w_i},$$

where

$$w_i = \frac{\phi_*(M_{*,i})}{d(M_{*,i})}$$

is a weight that corresponds to the inverse of the effective survey volume and  $\Delta \log M_{*,\text{bin}}$  is the bin width (typically 0.25). We also define

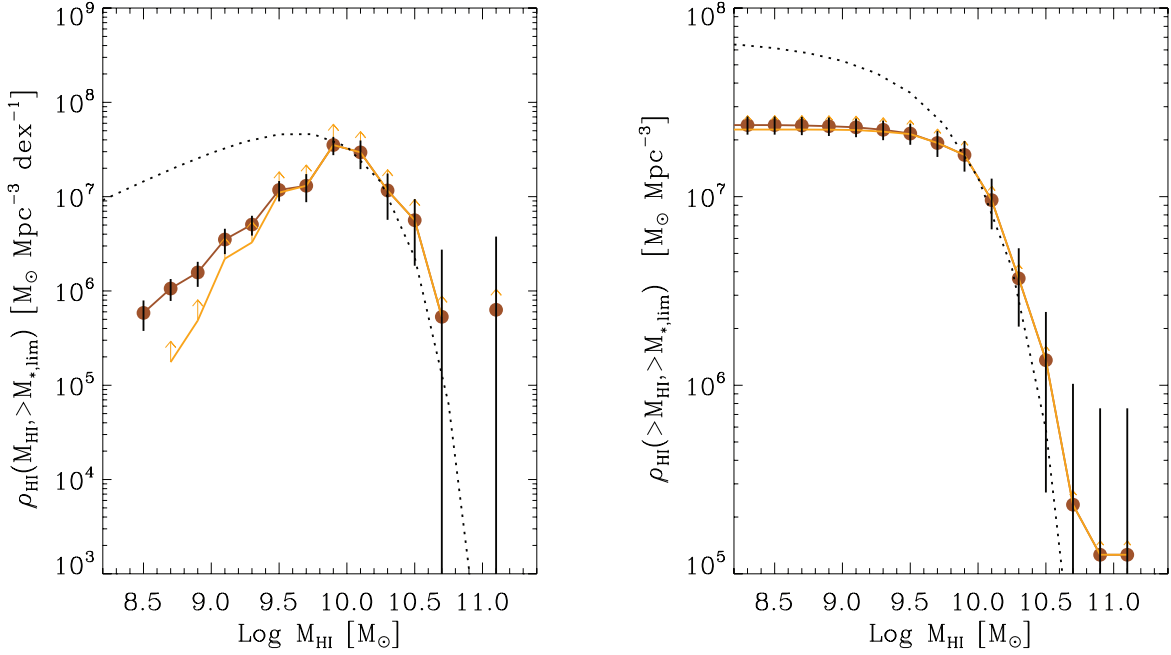
$$\bar{M}_{*,\text{bin}} = \frac{\sum_i M_{*,i} w_i}{\sum_i w_i}$$

as the weighted-average stellar mass of any given bin.

In the top panel of Fig. 2, we plot the local stellar mass function from Borch et al. (2006) and compare it with  $\phi_*(\bar{M}_{*,\text{bin}})$  based on the GASS DR1, which we find to be consistent with the input stellar mass function. We also plot the distribution of stellar masses in the (volume-limited) GASS PS which shows a deficit at high stellar masses, likely due to SDSS spectroscopic targeting criteria that avoid bright objects. Our method corrects for this incompleteness, which in any case is only significant for the highest stellar masses and is unlikely to impact our analysis.

We can use the same methodology to calculate  $\rho_{\text{SFR}}(M_*)$  and  $\rho_{\text{H I}}(M_*)$ , the volume densities of H I and SFR as a function of  $M_*$ . Again, because we are complete in any given  $M_*$  bin we find that

$$\rho_{\text{SFR}}(\bar{M}_{*,\text{bin}}) = \frac{\sum_i \text{SFR}_i \phi_*(M_{*,i}) w_i}{\sum_i w_i}$$



**Figure 3.** Derivation of cumulative H I cap mass found in massive galaxies ( $M_\star > 10^{10} M_\odot$ ). Yellow curve shows function derived using only H I detections and can be considered a lower limit. Brown curve and points show same function where we have included upper limits, which provides some indication of the upper bound. Dotted line shows total H I mass function derived from Zwaan et al. (2005). Left-hand panel: H I mass density versus H I mass. Right-hand panel: cumulative H I mass density above a given H I mass.

and

$$\rho_{\text{HI}}(\bar{M}_{\star, \text{bin}}) = \frac{\sum_i M_{\text{HI},i} \phi_\star(M_{\star,i}) w_i}{\sum_i w_i}.$$

Lastly, we apply a similar technique to determine the contribution of GASS galaxies to the overall H I mass density function  $\rho_{\text{HI}}(M_{\text{HI}})$  and SFR density function  $\rho_{\text{SFR}}(\text{SFR})$ . There are two caveats that should be noted with this calculation:

(i) The GASS sample will not necessarily be complete in any given  $M_{\text{HI}}$  or SFR bin, and therefore the derived function should only be taken as a lower bound on the total function.

(ii) If we perform the calculation by setting the H I mass for non-detections to the H I upper limit, we may misrepresent the shape of the density function below the point where the GASS sample is nearly complete. Because this calculation can provide an approximate upper bound on the actual function, we show these for illustrative purposes. However, these upper-limit mass functions should be treated with caution.

If we consider a complete bivariate distribution function, for example  $\phi(M_{\text{HI}}, M_\star)$ , we can evaluate a ‘partial’ distribution function

$$\phi(M_{\text{HI}}, > M_{\star, \text{lim}}) = \int_{M_{\star, \text{lim}}}^{\infty} \phi(M_{\text{HI}}, M_\star) dM_\star,$$

where we evaluate this function for a given stellar mass limit (e.g.  $M_{\star, \text{lim}} = 10^{10} M_\odot$  for the GASS). This partial distribution function can be simply calculated using our weight function derived above. Summing over all galaxies  $j$  where  $|\log M_{\text{HI},j} - \log M_{\text{HI}, \text{bin}}| < \Delta \log M_{\text{HI}, \text{bin}}/2$  and stellar mass  $M_{\star,j} > M_{\star, \text{lim}}$ , we find that

$$\phi(\bar{M}_{\text{HI}, \text{bin}}, > M_{\star, \text{lim}}) = \sum_j w_j.$$

As discussed above, we calculate this in two different ways to bracket our characterization of the function, either by omitting non-detections or by setting their H I mass to our calculated upper limit.

The partial H I mass density function is then

$$\rho(M_{\text{HI}}, > M_{\star, \text{lim}}) = M_{\text{HI}} \phi(M_{\text{HI}}, > M_{\star, \text{lim}})$$

and integrating, we obtain the cumulative function

$$\begin{aligned} \rho_{\text{HI}}(> M_{\text{HI}, \text{lim}}, > M_{\star, \text{lim}}) \\ = \int_{M_{\text{HI}, \text{lim}}}^{\infty} M'_{\text{HI}} \phi(M'_{\text{HI}}, > M_{\star, \text{lim}}) dM'_{\text{HI}}. \end{aligned}$$

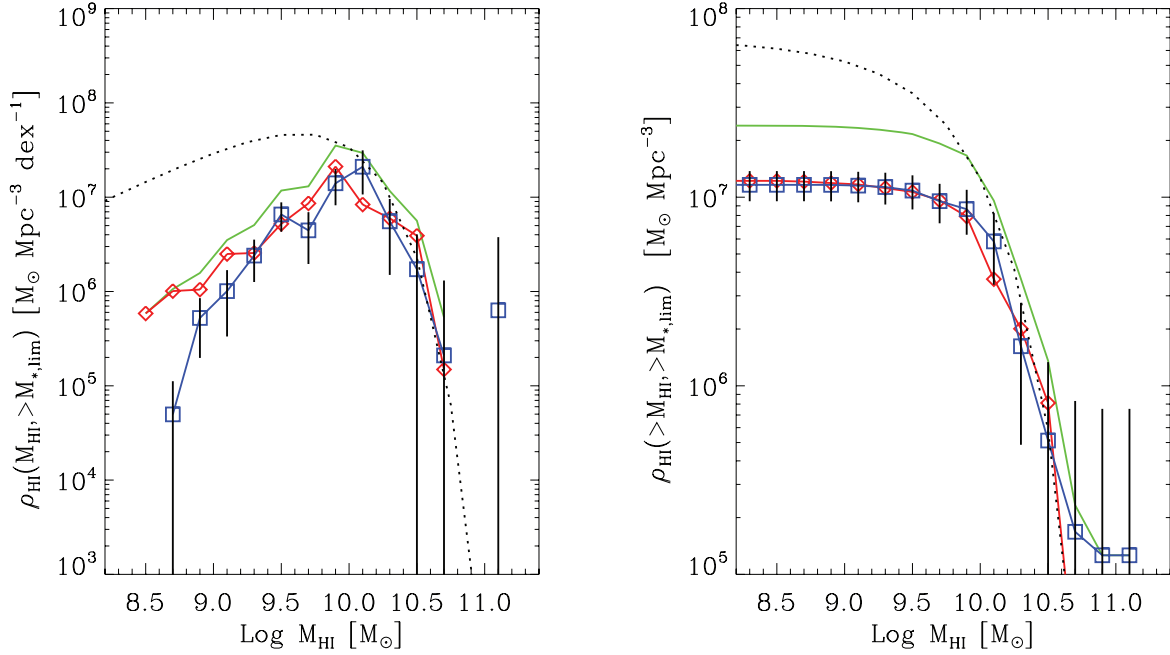
The integral of the cumulative H I mass density function provides an estimate of the total contribution of GASS galaxies to the H I density in the local Universe.

A similar calculation applies for  $\phi(\text{SFR}, M_\star)$ .

### 3.2 Volume-averaged H I, SFR, SFE

We show the partial H I mass density function calculated for the GASS sample in Fig. 3, and compare it to the total H I mass density distribution calculated using the Zwaan et al. (2005) H I mass function. In both plots, we show the H I mass density derived using only H I detections (yellow curve), which can be taken to be a lower limit, and the H I mass density using upper limits as discussed above (brown curve with points). Interestingly the GASS-derived curve matches the total H I mass density down to  $M_{\text{HI}} \sim 10^{10} M_\odot$ , suggesting that nearly all galaxies with high H I masses have stellar masses with  $M_\star > 10^{10} M_\odot$ . The GASS sample becomes increasingly less complete for H I masses below this limit.

We also find that GASS galaxies with  $\log M_{\text{HI}} > 9.75$  provide the bulk of the contribution to the total H I mass density, falling off significantly at lower gas masses. This can also be seen in the



**Figure 4.** Derivation of cumulative H I mass found in massive galaxies ( $\log M_{\star} > 10$ ) split by concentration. Left-hand panel: H I mass density versus H I mass. Right-hand panel: cumulative H I mass density above a given H I mass. Green histogram indicates full sample, blue histogram low-concentration ( $R_{90}/R_{50} < 2.6$ ) subsample, red histogram high concentration ( $R_{90}/R_{50} > 2.6$ ). Dotted line derived from Zwaan et al. (2005). Error bars shown only for low-concentration subsample.

sharp rise in the cumulative H I mass density  $\rho_{\text{HI}}(> M_{\text{HI}}, > M_{\star, \text{lim}})$  over this H I mass range. The cumulative mass density levels off at  $(2.5 \pm 0.3) \times 10^7 M_{\odot} \text{ Mpc}^{-3}$ ,  $36 \pm 5$  per cent of the total H I mass density derived using Zwaan et al. (2005). This integrated result shows that a significant fraction of the H I mass in the local Universe is associated with massive galaxies. There is little difference between the two curves, indicating that this result is robust with respect to the method that we use to account for non-detections.

We explore in Fig. 4 how this distribution changes if we split our sample according to a galaxy’s concentration, the ratio  $R_{90}/R_{50}$ , which we use as a proxy for distinguishing disc-dominated ( $R_{90}/R_{50} < 2.6$ ) and bulge-dominated ( $R_{90}/R_{50} > 2.6$ ) galaxies (as in e.g. Kauffmann et al. 2003). This cut yields 54 (133) low- (high-) concentration galaxies. We find that this split produces two nearly identical distributions, both contributing nearly equally to the total integrated H I mass density. For the GASS stellar mass range, disc-dominated galaxies do not account for the majority of the gas content of galaxies. Instead, H I appears to be evenly distributed across a range of galaxy types.

From our derived SFR density function, we measure a partial SFR density of  $(8.5 \pm 0.5) \times 10^{-3} M_{\odot} \text{ Mpc}^{-3} \text{ yr}^{-1}$  which is  $47 \pm 4$  per cent of the total SFR density (Salim et al. 2007; Wyder et al. 2007) and consistent with the results in those papers measured over the GASS stellar mass range. Interestingly this suggests that GASS galaxies account for nearly the same fraction of the total SFR density as they do in H I content. We can also calculate an integrated volume-averaged SFE and implied gas consumption time-scale which we determine to be  $2.9 \times 10^{-10} \text{ yr}^{-1}$  (time-scale =  $3.4 \pm 0.4$  Gyr) for GASS galaxies ( $M_{\star} > 10^{10} M_{\odot}$ ). Because we have measurements of the total H I and SFR density, we can calculate an SFE for lower mass galaxies ( $M_{\star} < 10^{10} M_{\odot}$ ), which we find to be  $2.1 \times 10^{-10} \text{ yr}^{-1}$  (time-scale =  $4.6 \pm 0.5$  Gyr). The volume-averaged SFE appears to be relatively constant across the full galaxy population.

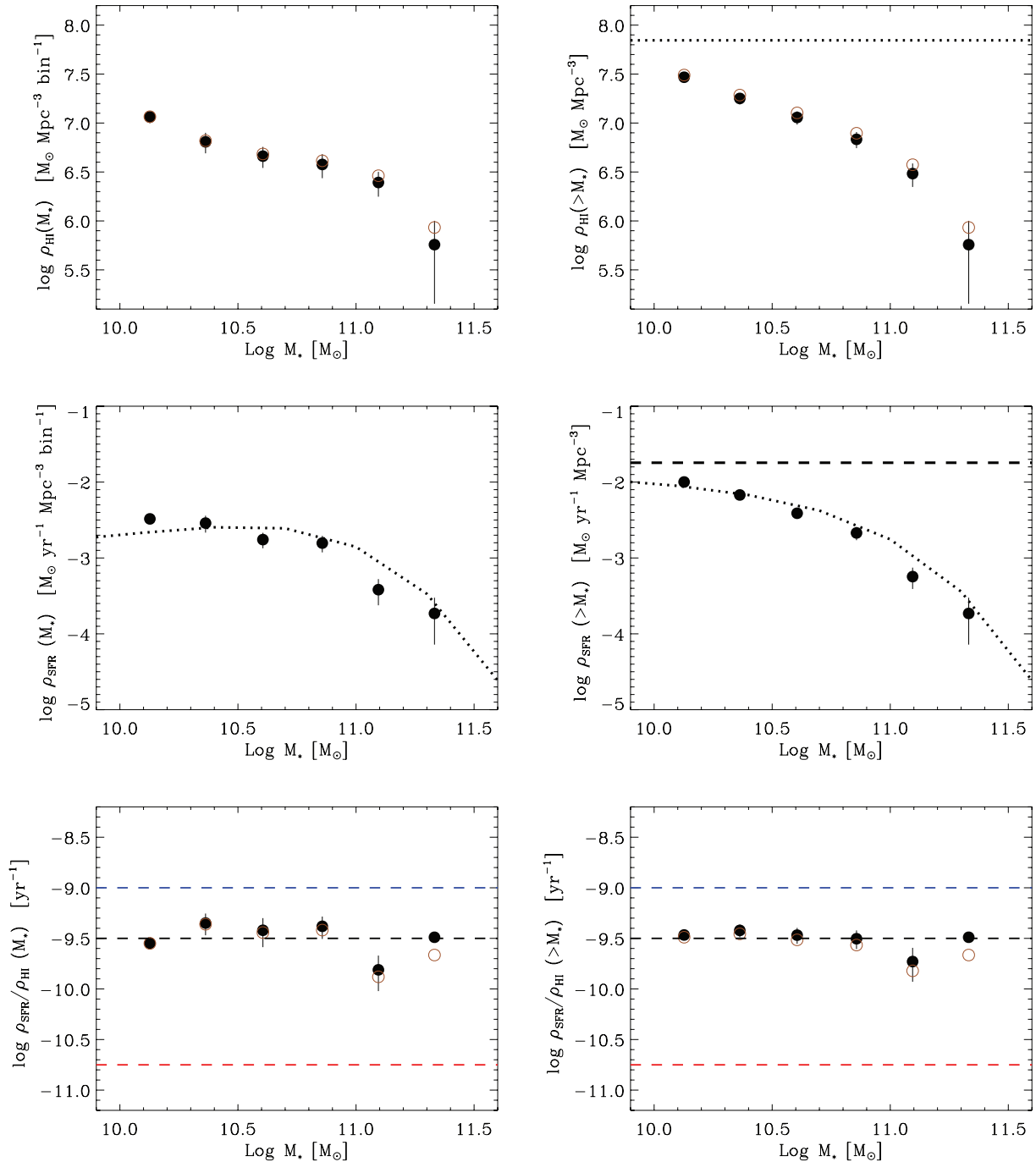
### 3.3 Global scaling relations: sSFR and SFE

Although our results suggest that the average SFE is similar above and below  $10^{10} M_{\odot}$ , it is possible that it only changes significantly above the transition mass at  $M_{\star} \sim 3 \times 10^{10} M_{\odot}$ . We show in Fig. 5 how the volume-averaged H I mass, SFR and SFE density functions vary with  $M_{\star}$ . Quantities are plotted per  $M_{\star}$  bin on the left, and as a cumulative quantity on the right. In the bottom panels, we see that the trend in gas consumption remains quite flat across all stellar masses, remaining constant up to our highest stellar mass bin, though with increasing scatter.

To explore this further, in Figs 6 and 7 we compare the global scaling relations of  $\text{SFR}/M_{\star}$  and SFE and find that they each paint a very different picture. In the left-hand panel of each pair of figures, we plot the binned sSFR ( $\Sigma \text{SFR}/\Sigma M_{\star}$ ) as a function of stellar mass, mass surface density, concentration and colour ( $M_{\star}$ ,  $\mu_{\star}$ ,  $R_{90}/R_{50}$ , and  $\text{NUV} - r$ ). In all cases,  $\text{SFR}/M_{\star}$  is steadily declining, typically by a factor of 10–30, across the range of the GASS sample. In the right-hand panels, we show the average SFE ( $\Sigma \text{SFR}/\Sigma M_{\text{HI}}$ , where H I non-detections are given zero H I mass). The average SFE is nearly flat, straddling  $\sim 3 \times 10^{-9} \text{ yr}^{-1}$ , corresponding to a gas depletion time-scale of 3 Gyr. Although in certain cases a small upward or downward trend is suggested, these deviations are not statistically significant. Values for these quantities are given in Tables 1 and 2. We have chosen to plot bin-averaged quantities where the numerator and denominator are summed separately in order to simplify the treatment of non-detections. We have checked that this trend is also apparent for mean or median  $\text{SFR}/M_{\star}$  and SFE measured for individual galaxies, the distribution of which is discussed in the next section.

This result is remarkable for two different reasons. The first is that the time-scale for gas consumption is nearly constant while the galaxy’s sSFR or ‘building time-scale’ is strongly dependent on stellar mass and correlated quantities.



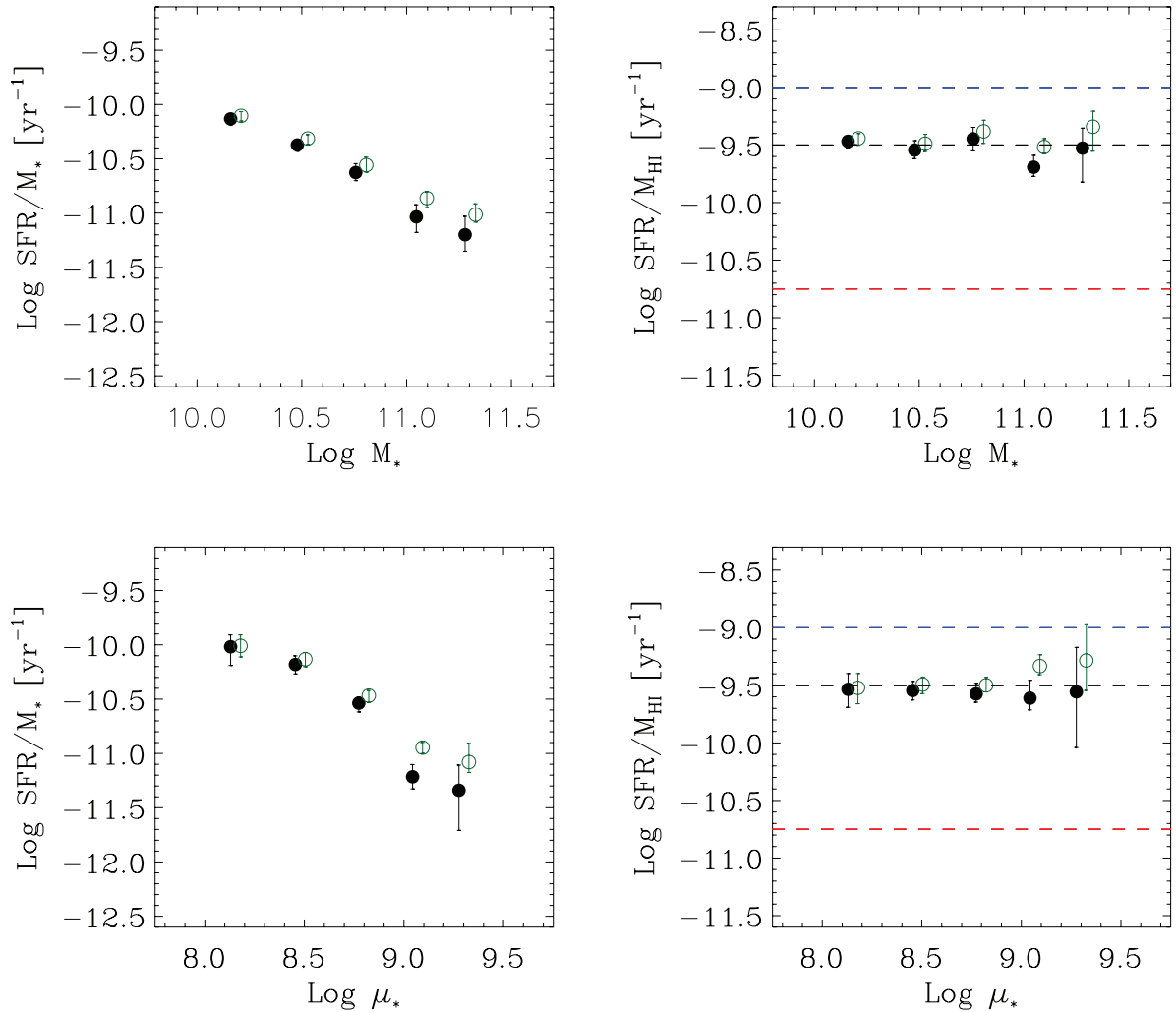


**Figure 5.** H I cap and SFR density and volume-averaged SFE versus  $\log M_*$ . For upper and lower panels, black dots show results derived from H I detections only, brown open dots are derived including non-detections, with H I masses set to calculated upper limits. Left-hand panels: densities shown per  $\log M_*$  bin (0.25 dex). Right-hand panels: cumulative density above a given  $M_*$ . Top panels: H I density,  $\rho_{\text{HI}}$ . Dotted line indicates cumulative H I density derived from Springob, Haynes & Giovanelli (2005). Middle panels: SFR density. Dotted lines show SFR density,  $\rho_{\text{SFR}}$  versus  $M_*$  from Schiminovich et al. (2007). Middle-right panel: dashed line indicates total value from Salim et al. (2007). Bottom panels: star formation efficiencies (or inverse gas consumption time-scales). Blue, black and red horizontal dashed lines correspond to high, average and low SFE values found for individual galaxies.

The second remarkable aspect is that the global average SFE that we measure for the GASS sample is very close to that observed locally for molecular gas in discs (e.g. Leroy et al. 2008, using the THINGS survey, Walter et al. 2008). This suggests that the entire H I reservoir is being converted into stars at the same rate as the molecular gas. It would appear to argue against a ‘bottleneck’ in

the flow of gas on to galaxies occurring, on average, at the interface between the atomic and molecular phases. Instead it would appear that the gas-limiting step occurs prior to the H I phase and that quenching of the detectable cold gas is not responsible for regulating the SF history of galaxies. We consider this result, and its relation to previous work, in the next section.





**Figure 6.** Mean sSFR (left-hand panels) and SFE (right-hand panels) as a function of derived quantities stellar mass,  $M_*$  (top panels) and stellar mass surface density  $\mu_*$  (bottom panels). Black points use SFRs calculated using dust correction  $A_{\text{NUV}}$  for ‘star-forming’ galaxies only (DC, D4NCUT sample), open points apply a dust correction over the full sample (DC sample). Error bars ( $1\sigma$ ) derived from bootstrap resampling. Blue, black and red horizontal dashed lines correspond to high, average and low SFE values found for individual galaxies. H I non-detections are given zero gas mass (see the text for details).

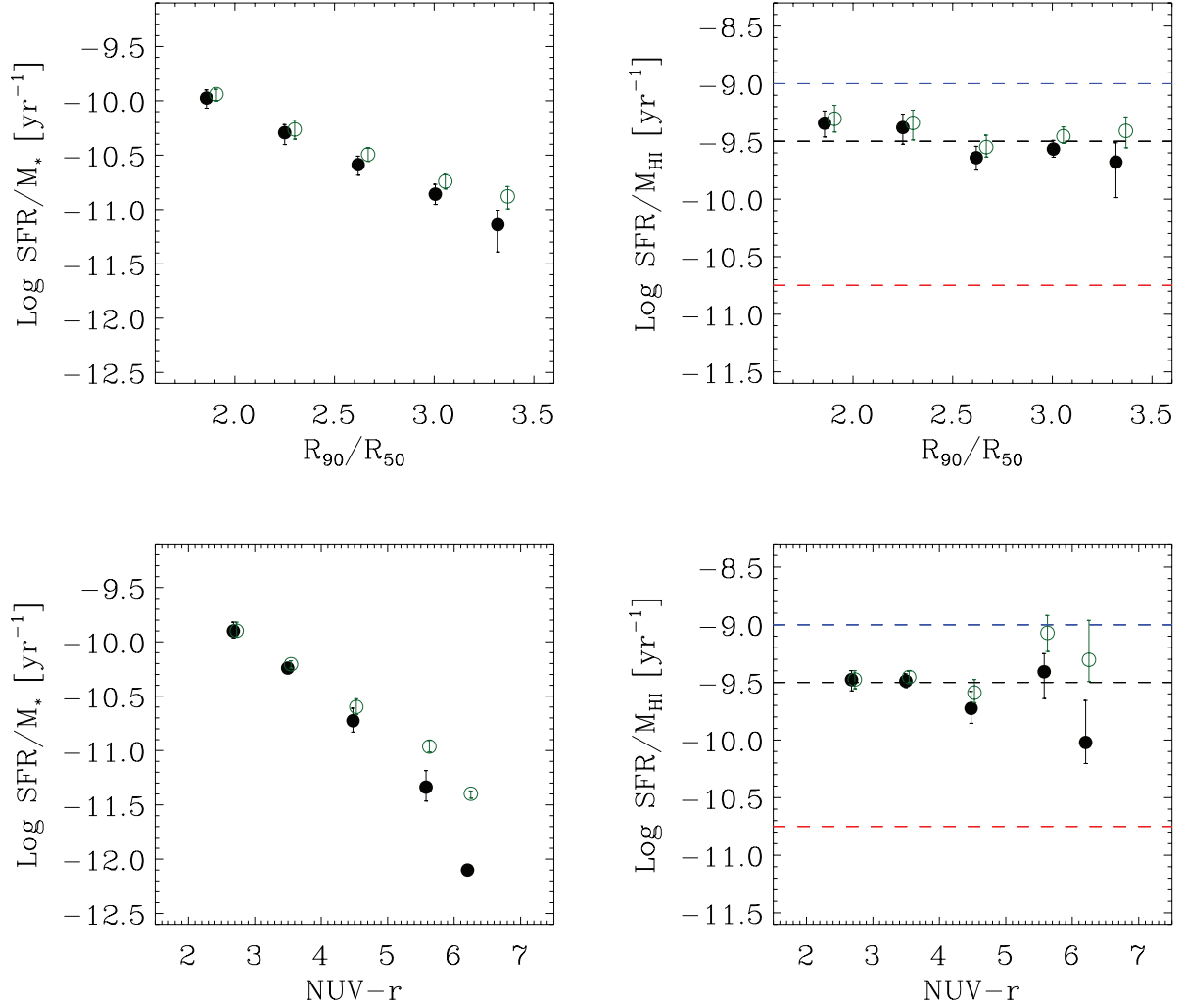
## 4 DISCUSSION

### 4.1 A nearly constant SFE

We first start out by asking whether a constant volume-averaged SFE is surprising, particularly for galaxies in the GASS mass range? While the global Schmidt law suggests an SFE that rises with gas surface density or varies inversely with dynamical or free-fall time-scale (Kennicutt 1998), more recent compilations of SF laws (e.g. Leroy et al. 2008) consider fixed giant molecular cloud star-forming efficiencies (with a varying atomic-to-molecular ratio) and pressure regulation of the SFE and/or atomic–molecular ratio. Because we measure neither the size of the gaseous and star-forming disc nor the molecular phase, it is not easy to connect our global results to these theoretical predictions. Global galaxy-averaged quantities sample a range of gas surface densities, time-scales and conditions within the interstellar medium. Additionally, our mean SFEs combine measurements from galaxies with range of morphological types, environments and presumably dark halo masses and spin parameters. It is a considerable theoretical challenge to interpret this result on its own.

Observationally, low efficiencies have been measured for low-mass galaxies (Geha et al. 2006), LSB galaxies (Boissier et al. 2008; Wyder et al. 2009) and DLAs (Wolfe & Chen 2006). Conversely high efficiencies have been measured in starburst galaxies (e.g. Lehnert & Heckman 1996; Kennicutt 1998) and in galaxies undergoing interactions (Young et al. 1986; Solomon & Sage 1988) and/or some form of environmental disturbance (e.g. stripping) (Koopmann & Kenney 2004; Rose et al. 2010). However, measurements of normal star-forming galaxies (e.g. Kennicutt 1998) show that galaxies follow the global Schmidt law, implying a slowly varying SFE over the range of gas surface densities typically probed.

Young et al. (1986) and Devereux & Young (1991) noted a flat SFE across the Hubble sequence using  $L(\text{FIR})/M(\text{H}_2)$  as a tracer. A similar result has been obtained by Boselli et al. (2001) and Boissier et al. (2001) who obtain a nearly constant SFE in a sample of normal spirals. More recently, Bothwell et al. (2009) suggested that the SFE is slowly increasing with galaxy luminosity (implying that more luminous galaxies have shorter gas consumption time-scales), but those data, which only include H I-detected galaxies, are nearly consistent with the constant values we have derived for the GASS sample.



**Figure 7.** Mean sSFR (left-hand panels) and SFE (right-hand panels) as a function of observed quantities concentration,  $R_{90}/R_{50}$  (top panels) and  $\text{NUV} - r$ , uncorrected for dust attenuation (bottom panels). Symbols and colours as in Fig. 6. H I non-detections are given zero gas mass (see the text for details).

#### 4.2 SFE versus $\text{SFR}/M_{\star}$

We return to the comparative question: why does the  $\text{SFR}/M_{\star}$  drop sharply with  $M_{\star}$  and  $\mu_{\star}$  while the SFE remains constant? Here we consider two different scenarios and leave a more detailed analysis for future work.

*Internal regulation at the atomic-to-molecular transition.* One possibility is that SF is inhibited within the gas reservoir, at the sink point rather than the supply location. Under the assumption that the efficiency of conversion of molecular gas into stars is nearly constant (e.g. Bigiel et al. 2008; Leroy et al. 2008), regulation would then occur at the interface between the atomic and molecular phases. Processes that can stabilize a gaseous disc, such as those proposed by Martig et al. (2009), are possible examples. As discussed in the Introduction, although such a process can explain a decreasing  $\text{SFR}/M_{\star}$  it may also result in a large reservoir of cold gas, leading to a decreasing SFE for galaxies that are actively being quenched. Therefore, it appears hard to reconcile internal regulation with the constant SFE and decreasing  $\text{SFR}/M_{\star}$  versus  $M_{\star}$  that we observe (instead one might expect some correlation between SFE and  $\text{SFR}/M_{\star}$ ). An alternative mechanism described in Hopkins, McClure-Griffiths & Gaensler (2008), with *in situ* formation of

molecular clouds in supershells, would also appear to be in conflict with our finding.

*Quenching or throttling of the H I supply.* A separate class of distinct processes are those that control the supply before (or quasi-instantaneously as) gas settles into the H I phase. This includes both the ‘ejective feedback’ and ‘preventive feedback’ quenching mechanisms discussed in Kereš et al. (2009). For example, SF (or AGN) may drive outflows, effectively ejecting gas from the system. AGN feedback may heat infalling gas, or prevent its cooling, and allow the build up of a reservoir of gas in a bound hidden phase (e.g. ionized, warm-hot). Although such a hidden phase has been invoked to balance accretion and SFRs in the Milky Way (Shull et al. 2009), more generally quenched gas in this phase may or may not provide an additional supply for SF through halo accretion, in that it may be permanently quenched or may subsequently cool on a longer time-scale (e.g. Dekel & Birnboim 2006). Additionally there are throttling scenarios that invoke varied gas accretion histories, essentially regulating the rate of infall on to the halo, with possible links to environmental conditions that lead to ‘starvation’ or ‘strangulation’ (Larson et al. 1980). Finally, this category also includes delayed or staged accretion histories (Boissier et al. 2001; Noeske et al. 2007) where the effective accretion time-scale varies

**Table 1.** Mean sSFRs within the GASS sample.

$x$	$\langle x \rangle$	$\text{Log}(\text{SFR}/M_\star \text{ yr})^a$	$\text{Log}(\text{SFR}/M_\star \text{ yr})^b$
$\text{Log } M_\star$	10.17	$-10.11 \pm 0.06$	$-10.08 \pm 0.05$
	10.49	$-10.47 \pm 0.11$	$-10.41 \pm 0.09$
	10.76	$-10.62 \pm 0.10$	$-10.55 \pm 0.08$
	11.05	$-11.08 \pm 0.09$	$-10.90 \pm 0.06$
	11.29	$-11.17 \pm 0.17$	$-11.01 \pm 0.07$
$\text{Log } \mu_\star$	8.13	$-9.92 \pm 0.16$	$-9.92 \pm 0.12$
	8.46	$-10.33 \pm 0.11$	$-10.25 \pm 0.08$
	8.77	$-10.61 \pm 0.07$	$-10.53 \pm 0.05$
	9.05	$-11.26 \pm 0.14$	$-10.98 \pm 0.07$
	9.27	$-11.38 \pm 0.33$	$-11.07 \pm 0.14$
$R_{90}/R_{50}$	1.85	$-10.07 \pm 0.12$	$-10.01 \pm 0.09$
	2.28	$-10.38 \pm 0.12$	$-10.34 \pm 0.12$
	2.61	$-10.63 \pm 0.08$	$-10.52 \pm 0.06$
	3.02	$-10.96 \pm 0.13$	$-10.81 \pm 0.09$
	3.32	$-11.45 \pm 0.22$	$-11.02 \pm 0.07$
$\text{NUV} - r$	2.71	$-9.93 \pm 0.09$	$-9.92 \pm 0.09$
	3.54	$-10.33 \pm 0.06$	$-10.29 \pm 0.05$
	4.50	$-10.83 \pm 0.10$	$-10.67 \pm 0.06$
	5.60	$-11.38 \pm 0.19$	$-11.00 \pm 0.07$
	6.18	$-12.09 \pm 0.02$	$-11.37 \pm 0.03$

<sup>a</sup>SFR is NUV-based with dust correction applied to galaxies with  $D_n(4000) < 1.7$  ( $\text{SFR}_{\text{NUV,DC,D4NCUT}}$ ). Uncertainties do not include possible systematic errors in SFR. <sup>b</sup>SFR is NUV-based with dust correction applied to all galaxies ( $\text{SFR}_{\text{NUV,DC}}$ ). Uncertainties do not include possible systematic errors in SFR.

**Table 2.** Mean SFEs of the GASS sample.

$x$	$\langle x \rangle$	$\text{Log}(\text{SFR}/M_{\text{H}} \text{ yr})^a$	$\text{Log}(\text{SFR}/M_{\text{H}} \text{ yr})^b$
$\text{Log } M_\star$	10.17	$-9.58 \pm 0.05$	$-9.54 \pm 0.04$
	10.49	$-9.52 \pm 0.10$	$-9.45 \pm 0.09$
	10.76	$-9.49 \pm 0.11$	$-9.41 \pm 0.11$
	11.05	$-9.74 \pm 0.10$	$-9.55 \pm 0.08$
	11.29	$-9.60 \pm 0.22$	$-9.46 \pm 0.13$
$\text{Log } \mu_\star$	8.13	$-9.64 \pm 0.13$	$-9.61 \pm 0.11$
	8.46	$-9.60 \pm 0.12$	$-9.54 \pm 0.11$
	8.77	$-9.63 \pm 0.07$	$-9.55 \pm 0.05$
	9.05	$-9.66 \pm 0.15$	$-9.38 \pm 0.11$
	9.27	$-9.50 \pm 0.43$	$-9.18 \pm 0.35$
$R_{90}/R_{50}$	1.85	$-9.48 \pm 0.07$	$-9.43 \pm 0.08$
	2.28	$-9.39 \pm 0.17$	$-9.35 \pm 0.15$
	2.61	$-9.68 \pm 0.12$	$-9.58 \pm 0.09$
	3.02	$-9.64 \pm 0.10$	$-9.49 \pm 0.08$
	3.32	$-9.94 \pm 0.24$	$-9.48 \pm 0.17$
$\text{NUV} - r$	2.71	$-9.51 \pm 0.10$	$-9.51 \pm 0.10$
	3.54	$-9.52 \pm 0.05$	$-9.46 \pm 0.05$
	4.50	$-9.83 \pm 0.13$	$-9.66 \pm 0.09$
	5.60	$-9.48 \pm 0.22$	$-9.09 \pm 0.12$
	6.18	$-9.95 \pm 0.16$	$-9.21 \pm 0.17$

<sup>a,b</sup>SFR and uncertainties as in Table 1.

with stellar mass. In principle, what all of these have in common is that they can produce significant variation in  $\text{SFR}/M_\star$  versus stellar mass while allowing the SFE to remain constant.

Interestingly, if interpreted in light of recent models explaining the SF law in atomic and molecular gas (Krumholz, McKee & Tumlinson 2008), our result suggests that on average most of the H I in GASS galaxies must reside at surface densities at or near the

atomic-to-molecular transition, with little or no systematic trend across the sample. Our conclusion, that quenching/throttling mechanisms appear to better explain our data, has already been hinted at in previous work, most notably Larson et al. (1980) and Boissier et al. (2001) both of which identified similar observational trends in disc galaxy samples. We note here that our average result does not preclude internal regulation from taking place in some, but not all, of the massive galaxies in our sample. It is possible that a ‘conspiracy’ resulting from galaxies with low, internally quenched SFE being balanced by galaxies with high SFEs could reproduce the mean scaling relation. We can investigate this by studying the actual distribution of SFE versus stellar mass and other properties across the sample, which we do below.

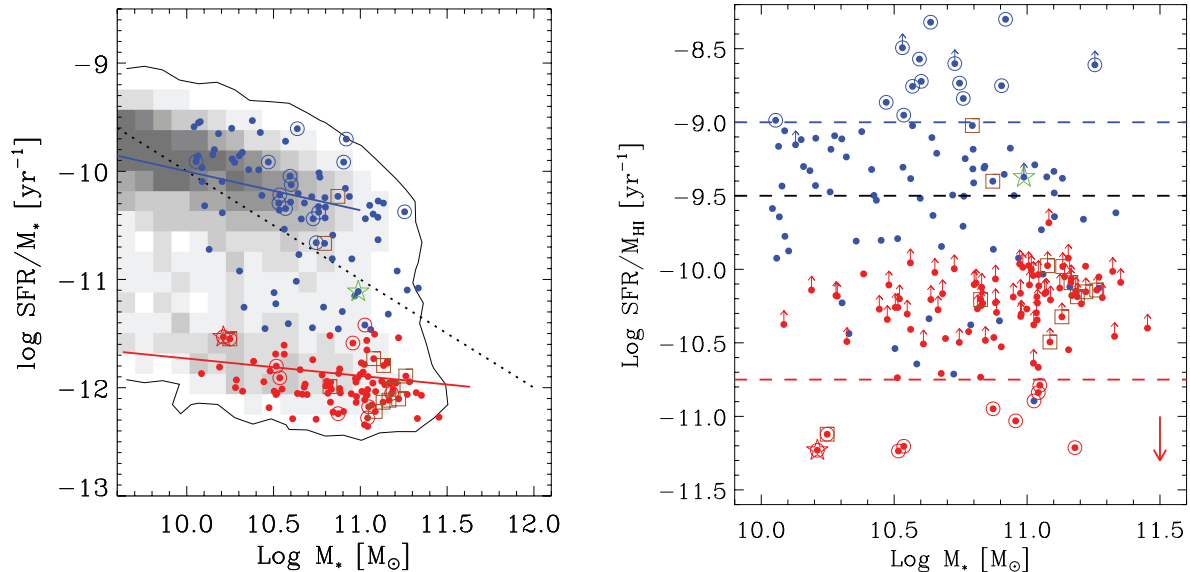
### 4.3 Distribution of SFE

Scatter around the mean relations discussed above may provide insight into the episodic nature of gas accretion and quenching in GASS galaxies, and we explore here the detailed distribution by plotting values and upper limits for individual galaxies in the DR1 sample. We note here that while large uncertainties in some of the derived SFRs and SFEs can obscure visible dependencies, such upper limits and sources of error for low-SFR or low-gas-mass galaxies are unlikely to affect the global trends discussed in previous sections. In general, their contribution to summed values is small. On the other hand, random and systematic errors will also shift high-SFR galaxies in these plots. As discussed above, we have verified that our global results are robust with respect to such possible sources of error, but errors in SFRs, typically  $\sim 0.3$  dex, may produce some of the outliers on the plots described here. Results quoted in this section should be treated as suggestive, motivating further study and improvements in the accuracy of these measures.

In Fig. 8, we plot the sSFR and SFEs as a function of stellar mass. Also shown is the distribution of  $\text{SFR}/M_\star$  for galaxies in the local Universe derived using volume-corrected UV-optical data from the *GALEX* and SDSS (Schiminovich et al. 2007). As expected, in any given stellar mass bin, GASS galaxies show a similar distribution in  $\text{SFR}/M_\star$  as the volume-corrected *GALEX*+SDSS sample, consistent with the fact that GASS galaxies are selected purely based on stellar mass, with no other selection bias. Our sample of massive galaxies is not purely passive; a significant fraction of GASS galaxies are on or close to the star-forming sequence. In fact, many of the galaxies are forming stars at rates higher than  $1 M_\odot \text{ yr}^{-1}$ . Among the population of massive galaxies are some objects that have some of the highest SFRs in the local Universe.

There is a locus of galaxies near  $\text{SFR}/M_\star \sim 10^{-12} \text{ yr}^{-1}$ , on the non-SF (red or dead) sequence. As discussed above and in Schiminovich et al. (2007), the  $\text{SFR}/M_\star$  for many of these galaxies is likely to be an overestimate due to the fact that UV light from evolved populations has not been subtracted when calculating SFRs. Although we do not label the red points as such, it is best to consider SFRs for these galaxies as an upper limit.

The right-hand side of Fig. 8 shows the distribution of SFEs across the sample. There is a broad spread of efficiencies in the sample with implied gas consumption time-scales ranging from less than 1 Gyr to over 100 Gyr. In fact, this spread is indeed much larger than that typically seen for molecular gas within disc galaxies (e.g. fig. 15 of Leroy et al. 2008). We highlight in this and future figures those galaxies with exceptionally high and low SFEs. Galaxies with  $\text{SFE} > 10^{-9} \text{ yr}^{-1}$  are forming stars at rates that are high compared to their present atomic gas mass (though still longer than likely dynamical time-scales, consistent with Lehnert



**Figure 8.** sSFR as a function of stellar mass (left-hand panel) and SFR efficiency as a function of stellar mass (right-hand panel). Grey-scale shows distribution of  $\text{SFR}/M_*$  for galaxies in the local Universe derived using volume-corrected UV-optical data from the *GALEX* and SDSS (Schiminovich et al. 2007). Blue/red points represent high/low sSFR, split at  $\log \text{SFR}/M_* = -11.5$ . Large blue circles are high-SFE galaxies ( $\text{SFE} > 10^{-9.0} \text{ yr}^{-1}$ ). Arrows indicate galaxies with no detection in H I. Large red circles are low-SFE galaxies ( $\text{SFE} < 10^{-10.75} \text{ yr}^{-1}$ ). Green star is gas-deficient object GASS 7050 and red star is gas-rich red galaxy GASS 3505 (see Paper I). Brown squares indicate objects with nearby companions that may also be detected within the Arecibo beam. In left-hand panel, the blue and red solid lines correspond to the star-forming sequence and the non-star-forming locus of red galaxies. A constant SFR of  $1 \text{ M}_\odot \text{ yr}^{-1}$  is shown by the black dotted line. In the right-hand panel, the blue, black and red lines are used to show approximate locations of high, average and low SFE. Red arrow in the bottom right corner denotes amount SFR values of red points would move if 50 per cent of UV light came from evolved stars.

& Heckman 1996). Galaxies with  $\text{SFE} < 10^{-10.75} \text{ yr}^{-1}$  appear to be passive in comparison to their current gas content. GASS H I non-detections are shown with upward arrows, suggesting that their efficiencies could be higher. The vast majority of the non-detections also have very low sSFRs which could suggest that estimates of their SFRs and SFEs are too high. As a result, the red points that are H I non-detections may have higher *or* lower SFEs, and are poorly constrained in this diagram. Overall, the broad scatter in SFEs versus stellar mass does not point to a strong correlation between the two.

The high-SFE galaxies, indicated with large blue circles, are the galaxies with the shortest gas consumption time-scales. These galaxies are predominantly above the transition mass,  $M_* \sim 3 \times 10^{10} \text{ M}_\odot$ . Although these galaxies have elevated SFEs, they do not occupy an unusual location in the  $\text{SFR}/M_*$  versus  $M_*$  diagram, with most situated on or near the star-forming sequence. This suggests that the reason why their SFE is high is because they have low gas mass for their given present averaged SFR.

Stellar mass surface density is known to display a strong correlation with sSFR (Kauffmann et al. 2003), and we show that the GASS sample displays a similar correlation in Fig. 9. Above a transition stellar mass surface density of  $\mu_* > 10^{8.5} \text{ M}_\odot \text{ kpc}^{-2}$  a much greater spread of  $\text{SFR}/M_*$  is observed at a given  $\mu_*$ . This transition threshold has already been noted in Paper I and in previous work (Kauffmann et al. 2003; Brinchmann et al. 2004). On the other hand, the SFE, unlike the trend versus stellar mass, shows a rather curious ‘T-shaped’ distribution. Below the  $\mu_*$  transition, the SFEs are nearly constant, straddling the line around  $\text{SFE} = 10^{-9.5} \text{ yr}^{-1}$ . Above the transition, the distribution in SFE becomes extremely broad, with very high and very low efficiencies present. The highest SFE galaxies occupy the ‘knee’ of the  $\text{SFR}/M_*$  versus  $\mu_*$  diagram and might themselves represent a transition population (see e.g. Salim et al. 2007). The low-SFE galaxies, on the other hand, appear to reside with passive galaxies. These galaxies host a

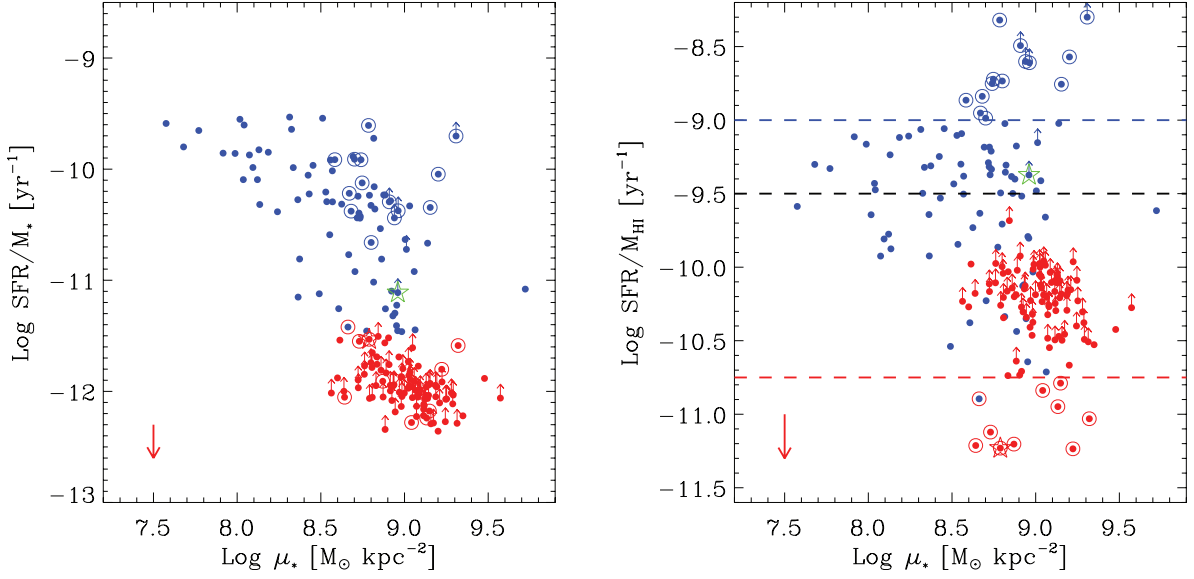
sufficiently large gaseous reservoir that subsequent future SF may lead to evolution off of the red sequence. Our results highlight the fact that the combination of gas, SF and stellar mass measurements can be useful for obtaining information about the possible future evolution of massive galaxies. We explore this further in Fig. 10 by investigating three plots that link gas fraction to  $\text{SFR}/M_*$  and SFE.

In the first plot comparing SFE versus H I gas mass fraction (upper left), we see that high-SFE galaxies have low gas fractions, with the converse being true for low-SFE galaxies. Star-forming and green valley galaxies display a range of gas fractions from 2 to 100 per cent. We plot SFE versus  $\text{SFR}/M_*$  (upper right) and find that the highest SFE galaxies do not have the highest sSFRs, but instead peak at slightly lower values. This suggests that a high SFE is not driven by a high SFR. Not surprisingly, the gas fraction versus  $\text{SFR}/M_*$  relation (lower right) shows that star-forming and green valley galaxies reveal correlated gas fraction and sSFRs with the extreme SFE galaxies lying off of this relation. All of these plots, taken together, suggest that it is low H I content, as opposed to an excess in SFR, that is responsible for the majority of high-SFE galaxies.

Figs 8–10 can be used to isolate galaxies that show gas excess or deficiency when compared to their current SFRs, useful for determining what causes high or low SFE in galaxies. Fig. 11 attempts to combine this diagnostic capability in one diagram, by plotting the offset of a galaxy’s SFR and H I mass relative to the average value for a given stellar mass.

We use data from this paper and Paper I and perform linear regression fits to the mean SFR and  $M_{\text{HI}}$  versus  $M_*$ .<sup>2</sup> For the H I masses

<sup>2</sup> We note that we have used the average SFR versus  $M_*$  as opposed to the relationship between SFR and  $M_*$  along the star-forming sequence described in Schiminovich et al. (2007).



**Figure 9.** sSFR (left-hand panel) and SFR efficiency (right-hand panel) as a function of stellar mass surface density. Colours and symbols as in Fig. 8.

we use the mean values for which the masses of non-detections have been set to the upper limit (Paper I, table 4, column 1). We find that

$$\log < M_{\text{HI}}(M_*) > = 0.02 \log M_* + 9.52$$

and

$$\log < \text{SFR}(M_*) > = 0.15 \log M_* - 1.5$$

and define

$$\Delta \log \frac{M_{\text{HI}}}{M_*} = [\log M_{\text{HI}} - \log < M_{\text{HI}}(M_*) >]$$

$$\Delta \log \frac{\text{SFR}}{M_*} = [\log \text{SFR} - \log < \text{SFR}(M_*) >]$$

In this delta gas fraction versus delta sSFR plot, four quadrants are delineated. The top-right quadrant contains those galaxies with high  $\text{SFR}/M_*$  and  $M_{\text{HI}}/M_*$ , in other words galaxies that are gas-rich and actively star-forming. The majority of the GASS detections fall here along a line where gas fraction excess equals the  $\text{SFR}/M_*$  excess which corresponds approximately to a line of constant SFE. This quadrant is also likely to include gas-rich mergers, starbursts and high surface brightness galaxies with extended UV discs (Thilker et al. 2007). The upper-left quadrant contains gas-rich galaxies with lower than average SFRs, including the low-SFE galaxies highlighted above. Low surface brightness galaxies and galaxies that have recently accreted gas might be found here. The majority of galaxies identified as low SFE are forming stars below average rates, consistent with their being passive galaxies with large reservoirs of gas. These galaxies do not appear to have extremely high gas masses and average SFRs (which may be more typical of galaxies with lower stellar mass). GASS 3505, an unusually gas-rich galaxy discussed in Paper I, belongs to this class of object and is indicated with a red star in the figure.

The lower-left quadrant contains gas-poor galaxies with below average SFR. Passive galaxies are found in this quadrant. Finally the lower-right quadrant contains high-SFE galaxies that are relatively gas poor but with above average SFR. On this plot, it becomes apparent that these galaxies are distributed across the quadrant and appear to have either high SFRs when compared to their (typical)

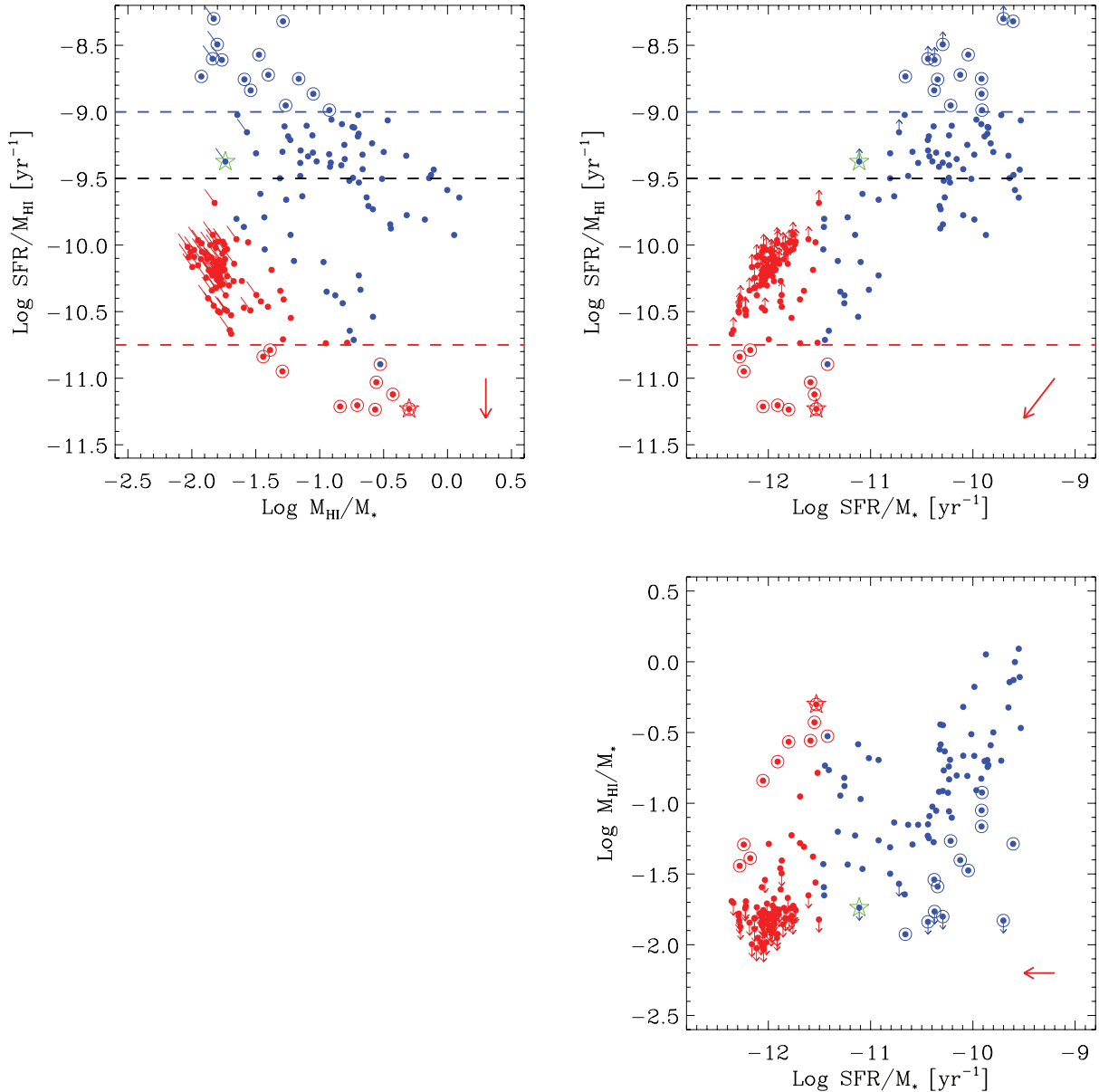
gas mass or lower than average gas mass but typical SFRs. The most gas deficient in this category has possibly very recently experienced a process that disrupted gas flow and/or removed gas, such as starvation or stripping. The latter scenario is likely to produce gas deficiencies even lower than we have observed. Later stages of such galaxies might also be found in the third quadrant below or near the line of constant SFE. GASS 7050, identified in Paper I as an unusual gas-poor galaxy, appears to belong to this category and is indicated with a green star in the figure.

Lastly, we return to the question of whether or not the scatter around our mean scaling relations suggests a diversity in the processes that trigger and quench SF in the GASS sample. Without question the GASS sample does not show a tight distribution in SFEs, and it is tempting to reconsider internal quenching mechanisms for at least some fraction of systems. Alternately, our results may rule out scenarios where the inflow of material is occurring as a steady flow (or drizzle) on to the galaxy. Instead, the scatter may be an indication that accretion of H I is episodic, with infalling gas arriving in larger discrete chunks. It may suggest that large-scale processes play a role in regulating the growth and evolution of gas and SF in galaxies.

Future work is being planned to investigate the large- and small-scale environment of GASS galaxies, and in particular the outliers, which should reveal whether environmental processes are driving this evolution. Some galaxies with high SFRs may be undergoing a merger or interaction that is driving up the SFR in these galaxies, and it will also be interesting to investigate the connection between signs of interaction, merging or other disturbance and location on this diagram.

## 5 CONCLUSIONS

We use measurements of the H I content, stellar mass and SFRs in  $\sim 190$  massive galaxies with  $M_* > 10^{10} M_\odot$ , obtained from the GALEX Arecibo SDSS survey described in Paper I to explore the global scaling relations associated with the ratio  $\text{SFR}/M_{\text{HI}}$ , which we call the H I-based SFE. We find that:



**Figure 10.** Projected views of the three-space of SFE versus gas fraction versus sSFR. SFE versus  $M_{\text{HI}}/M_{\star}$  (top-left panel) and versus  $\text{SFR}/M_{\star}$  (top-right panel). Lower panel shows  $M_{\text{HI}}/M_{\star}$  as a function of  $\text{SFR}/M_{\star}$ . Colours and symbols as in Fig. 8. In the top-left panel, H I masses are used to calculate both axes – upper limits are indicated by the short diagonal lines.

(i) We can measure the volume-averaged H I mass density and SFR density for the GASS sample. GASS galaxies account for  $36 \pm 5$  per cent of the total H I mass density and  $47 \pm 5$  per cent of the SFR mass density.

(ii) Unlike the mean sSFR, which decreases with stellar mass and stellar mass surface density, the SFE remains relatively constant across the sample with a value close to  $\text{SFE} = 10^{-9.5} \text{ yr}^{-1}$  (or an equivalent gas consumption time-scale of  $\sim 3 \times 10^9 \text{ yr}$ ). We find little variation in SFE with stellar mass, stellar mass surface density,  $\text{NUV} - r$  colour and concentration. We interpret these results as an indication that external processes or feedback mechanisms that control the gas supply are important for regulating SF in massive galaxies.

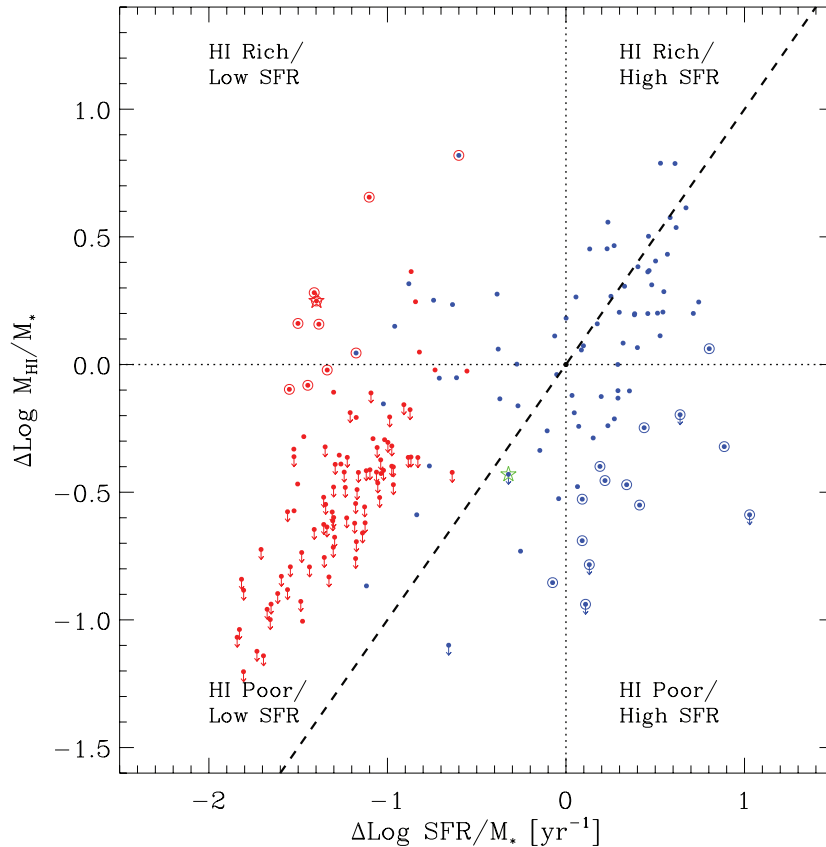
(iii) Approximately 5 per cent of the sample shows high efficiencies with  $\text{SFE} > 10^{-9} \text{ yr}^{-1}$ , and we suggest that this is very

likely due to a deficiency of cold gas rather than an excess SFR. Conversely, we also find a similar fraction of galaxies that appear to be gas-rich for their given sSFR, although such galaxies show both a higher than average gas fraction and lower than average sSFR. Both of these populations are plausible candidates for ‘transition’ galaxies, with potential for a change (either decrease or increase) in their sSFR in the near future.

## ACKNOWLEDGMENTS

We thank the Arecibo staff, in particular Phil Perillat, Ganesan Rajagopalan and the telescope operators for their assistance, and Hector Hernandez for scheduling the observations. RG and MPH acknowledge support from NSF grant AST-0607007 and from the Brinson Foundation.





**Figure 11.** Gas fraction excess versus sSFR excess. Colours and symbols as in Fig. 8. Zero-delta lines are plotted, as well as the line indicating  $\Delta \text{Log SFR}/M_* = \Delta \text{Log } M_{\text{HI}}/M_*$ .

The Arecibo Observatory is part of the National Astronomy and Ionosphere Center, which is operated by Cornell University under a cooperative agreement with the National Science Foundation.

*GALEX* is a NASA Small Explorer, launched in 2003 April. We gratefully acknowledge NASA's support for construction, operation and science analysis for the *GALEX* mission, developed in cooperation with the Centre National d'Etudes Spatiales of France and the Korean Ministry of Science and Technology.

Funding for the SDSS and SDSS-II has been provided by the Alfred P. Sloan Foundation, the participating institutions, the National Science Foundation, the US Department of Energy, the NASA, the Japanese Monbukagakusho, the Max Planck Society and the Higher Education Funding Council for England. The SDSS Web Site is <http://www.sdss.org/>.

The SDSS is managed by the Astrophysical Research Consortium for the participating institutions. The participating institutions are the American Museum of Natural History, Astrophysical Institute Potsdam, University of Basel, University of Cambridge, Case Western Reserve University, University of Chicago, Drexel University, Fermilab, the Institute for Advanced Study, the Japan Participation Group, Johns Hopkins University, the Joint Institute for Nuclear Astrophysics, the Kavli Institute for Particle Astrophysics and Cosmology, the Korean Scientist Group, the Chinese Academy of Sciences (LAMOST), Los Alamos National Laboratory, the Max Planck Institute for Astronomy, the MPA, New Mexico State University, Ohio State University, University of Pittsburgh, University

of Portsmouth, Princeton University, the United States Naval Observatory and the University of Washington.

## REFERENCES

- Bell E. F., McIntosh D. H., Katz N., Weinberg M. D., 2003, *ApJS*, 149, 289
- Bigiel F., Leroy A., Walter F., Brinks E., de Blok W. J. G., Madore B., Thornley M. D., 2008, *AJ*, 136, 2846
- Boissier S., Boselli A., Prantzos N., Gavazzi G., 2001, *MNRAS*, 321, 733
- Boissier S. et al., 2008, *ApJ*, 681, 244
- Boselli A., Gavazzi G., Donas J., Scodreggio M., 2001, *AJ*, 121, 753
- Borch A. et al., 2006, *A&A*, 453, 869
- Bothwell M. S., Kennicutt R. C., Lee J. C., 2009, *MNRAS*, 400, 154
- Bregman J. N., Hogg D. E., Roberts M. S., 1992, *ApJ*, 387, 484
- Brinchmann J., Charlot S., White S. D. M., Tremonti C., Kauffmann G., Heckman T., Brinkmann J., 2004, *MNRAS*, 351, 1151
- Calzetti D., Armus L., Bohlin R. C., Kinney A. L., Koornneef J., Storchi-Bergmann T., 2000, *ApJ*, 533, 682
- Catinella B. et al., 2010, *MNRAS*, 403, 683 (Paper I)
- Chabrier G., 2003, *PASP*, 115, 763
- Cortese L., Boselli A., Franzetti P., Decarli R., Gavazzi G., Boissier S., Buat V., 2008, *MNRAS*, 386, 1157
- Dekel A., Birnboim Y., 2006, *MNRAS*, 368, 2
- Devereux N. A., Young J. S., 1991, *ApJ*, 371, 515
- Donas J. et al., 2007, *ApJS*, 173, 597
- Geha M., Blanton M. R., Masjedi M., West A. A., 2006, *ApJ*, 653, 240
- Giovanelli R. et al., 2005, *AJ*, 130, 2598
- Hanish D. J. et al., 2006, *ApJ*, 649, 150



- Hopkins A. M., McClure-Griffiths N. M., Gaensler B. M., 2008, *ApJ*, 682, L13
- Johnson B. D. et al., 2006, *ApJ*, 644, L109
- Johnson B. D. et al., 2007, *ApJS*, 173, 377
- Kauffmann G. et al., 2003, *MNRAS*, 341, 33
- Kennicutt R. C., Jr, 1983, *ApJ*, 272, 54
- Kennicutt R. C., Jr, 1998, *ApJ*, 498, 541
- Kennicutt R. C., Jr, Tamblyn P., Congdon C. E., 1994, *ApJ*, 435, 22
- Kennicutt R. C. et al., 2009, *ApJ*, 703, 1672
- Kereš D., Katz N., Fardal M., Davé R., Weinberg D. H., 2009, *MNRAS*, 395, 160
- Koopmann R. A., Kenney J. D. P., 2004, *ApJ*, 613, 851
- Krumholz M. R., McKee C. F., Tumlinson J., 2008, *ApJ*, 689, 865
- Larson R. B., Tinsley B. M., Caldwell C. N., 1980, *ApJ*, 237, 692
- Lehnert M. D., Heckman T. M., 1996, *ApJ*, 472, 546
- Leroy A. K., Walter F., Brinks E., Bigiel F., de Blok W. J. G., Madore B., Thornley M. D., 2008, *AJ*, 136, 2782
- Martig M., Bournaud F., Teyssier R., Dekel A., 2009, *ApJ*, 707, 250
- Martin A. M., Giovanelli R., Haynes M. P., Saintonge A., Hoffman G. L., Kent B. R., Stierwalt S., 2009, *ApJS*, 183, 214
- Meurer G. R. et al., 2006, *ApJS*, 165, 307
- Meyer M. J. et al., 2004, *MNRAS*, 350, 1195
- Morganti R. et al., 2006, *MNRAS*, 371, 157
- Noeske K. G. et al., 2007, *ApJ*, 660, L47
- Oosterloo T. A., Morganti R., Sadler E. M., van der Hulst T., Serra P., 2007, *A&A*, 465, 787
- Rich R. M. et al., 2005, *ApJ*, 619, L107
- Roberts M. S., 1963, *ARA&A*, 1, 149
- Rose J. A., Robertson P., Miner J., Levy L., 2010, *AJ*, 139, 765
- Saintonge A., Giovanelli R., Haynes M. P., Hoffman G. L., Kent B. R., Martin A. M., Stierwalt S., Brosch N., 2008, *AJ*, 135, 588
- Salim S. et al., 2007, *ApJS*, 173, 267
- Schiminovich D. et al., 2007, *ApJS*, 173, 315
- Shull J. M., Jones J. R., Danforth C. W., Collins J. A., 2009, *ApJ*, 699, 754
- Solomon P. M., Sage L. J., 1988, *ApJ*, 334, 613
- Springob C. M., Haynes M. P., Giovanelli R., 2005, *ApJ*, 621, 215
- Stierwalt S., Haynes M. P., Giovanelli R., Kent B. R., Martin A. M., Saintonge A., Karachentsev I. D., Karachentseva V. E., 2009, *AJ*, 138, 338
- Thilker D. A. et al., 2007, *ApJS*, 173, 538
- Walter F., Brinks E., de Blok W. J. G., Bigiel F., Kennicutt R. C., Thornley M. D., Leroy A., 2008, *AJ*, 136, 2563
- Wang L., Rowan-Robinson M., 2009, *MNRAS*, 398, 109
- Wolfe A. M., Chen H. W., 2006, *ApJ*, 652, 981
- Wyder T. K. et al., 2007, *ApJS*, 173, 293
- Wyder T. K. et al., 2009, *ApJ*, 696, 1834
- Yi S. K. et al., 2005, *ApJ*, 619, L111
- Young J. S., Kenney J. D., Tacconi L., Claussen M. J., Huang Y. L., Tacconi-Garman L., Xie S., Schloerb F. P., 1986, *ApJ*, 311, L17
- Zwaan M. A., Meyer M. J., Staveley-Smith L., Webster R. L., 2005, *MNRAS*, 359, L30

This paper has been typeset from a  $\text{\LaTeX}$  file prepared by the author.